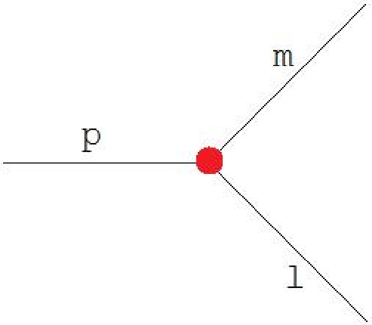
Proton decay matrix elements from chirally symmetric lattice QCD

▷ Paul Cooney, The University of Edinburgh, RBC-UKQCD collaboration
▷ Workshop on Underground Detectors Investigating Grand Unification (UDiG) at Brookhaven National Laboratory

$$p \to m + l$$



For a generic decay channel, the partial decay width is:

$$\Gamma(p \to m{+}\bar{l}) = \left[\frac{m_p}{32\pi^2}\left(1 - \left(\frac{m_m}{m_p}\right)^2\right)^2\right] \left|\sum_i C^i W_0^i(p \to m + \bar{l})\right|^2$$

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The form factors can be related to a matrix element

$$P_L \left[W_0^i(q^2) - i\not{q} W_q^i(q^2) \right] u(k, s) = \langle m | \mathcal{O}^i | N \rangle$$

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The operators $\mathcal{O}^i$ are given by

$$\begin{aligned}\mathcal{O}^{RL} &= \epsilon^{abc} u^{a,T}(x,t) C P_R d^b(x,t) P_L u^c(x,t)\\ \mathcal{O}^{LL} &= \epsilon^{abc} u^{a,T}(x,t) C P_L d^b(x,t) P_L u^c(x,t)\end{aligned}$$

Define a general operator of the form

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where Γ_i are matrices with two spin indices, labelled by,

$$\begin{array}{cc} S = 1 & P = \gamma_5 \\ V = \gamma_\mu & A_\mu = \gamma_\mu \gamma_5 \\ T = \frac{1}{2}\{\gamma_\mu, \gamma_\nu\} & \tilde{T} = \gamma_5 \frac{1}{2}\{\gamma_\mu, \gamma_\nu\} \\ R = P_R = \frac{1}{2}(1+\gamma_5) & L = P_L = \frac{1}{2}(1-\gamma_5) \end{array}$$

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Operators with this structure are also used later in nucleon correlation functions and in the non-perturbative renormalization

We could measure the matrix elements $\langle m|\mathcal{O}^i|N\rangle$ directly

- Known as the *direct* method
- Three-point functions are required
- Computationally expensive

Alternatively can relate the three-point functions to two-point functions using Chiral Perturbation Theory

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- Computationally cheaper
- Introduces an additional source of error

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For $p \rightarrow \pi^0 + e^+$, the chiral perturbation theory gives

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W_0^{RL}(p \rightarrow \pi^0 + e^+) &= \alpha(1 + D + F)/\sqrt{2}f + \mathcal{O}(m_l^2/m_N^2) \\
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α and β are low energy constants from the chiral lagrangian
They can be calculated from two–point functions

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Define a class of two–point functions

$$f_{\Gamma_1\Gamma_2,\Gamma_3\Gamma_4}(t) = \sum_x \mathrm{tr}\left[\langle \mathcal{O}^{\Gamma_1\Gamma_2}\bar{\mathcal{O}}^{\Gamma_3\Gamma_4}\rangle P\right]$$

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Example: the proton correlation function

$$\sum_x \langle J_p(x,t)\bar{J}_p(0)\rangle = f_{PS,PS}(t)$$

Strategy:

- First find m_N from a correlated fit to the effective mass

$$m_{\text{eff}}(t) = \log\left(\frac{f_{PS,PS}(t)}{f_{PS,PS}(t+1)}\right) \to m_N \quad t \gg 0$$

- Then find G_N from a correlated fit to an effective amplitude

$$G_{N,\text{eff}} = \sqrt{2 f_{PS,PS}\, \mathrm{e}^{m_N t}} \to G_N \quad t \gg 0$$

- Finally to calculate α and β we use a ratio of two–point functions

$$R_\alpha(t) = 2G_N \frac{f_{RL,PS}(t)}{f_{PS,PS}(t)} \to \alpha \quad R_\beta(t) = 2G_N \frac{f_{LL,PS}(t)}{f_{PS,PS}(t)} \to \beta$$

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  - Nearly exact chiral symmetry
  - Inverse lattice spacing $a^{-1} = 1.73(3)$ GeV
- Two different lattice volumes
  - $V = 16^3 \times 32 \approx 1.8\text{fm}^3$
  - $V = 24^3 \times 64 \approx 2.7\text{fm}^3$
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Fitting

Fit by minimising a correlated χ^2

$$\chi^2(p) = \sum_{t,t'} \left[p_{\text{eff}}(t) - p\right] C^{-1}_{tt'} \left[p_{\text{eff}}(t') - p\right]$$

With correlation Matrix

$$C_{tt'} = \frac{1}{N_{\text{boot}}} \sum_{n=1}^{N_{\text{boot}}} \left[p^{(n)}_{\text{eff}}(t) - \bar{p}_{\text{eff}}(t)\right] \left[p^{(n)}_{\text{eff}}(t') - \bar{p}_{\text{eff}}(t')\right].$$

Bootstrap to get central value and errors

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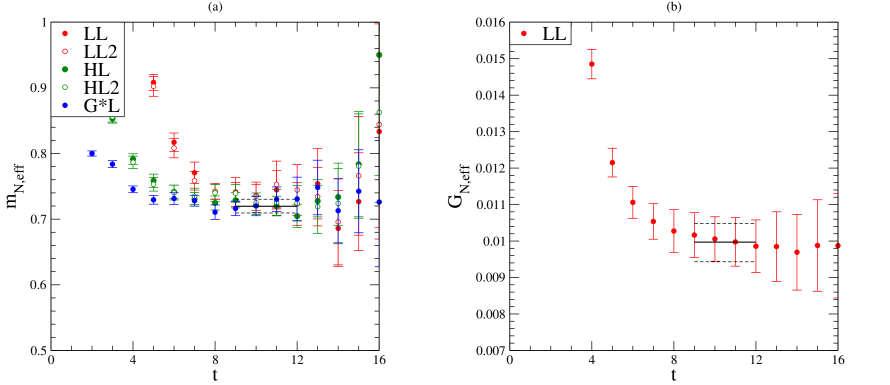
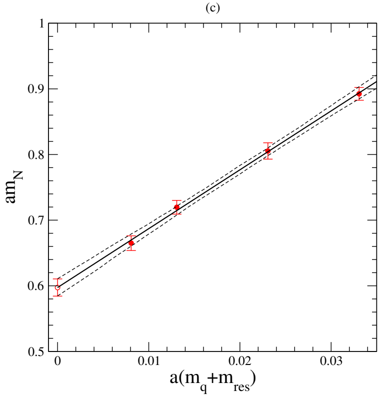
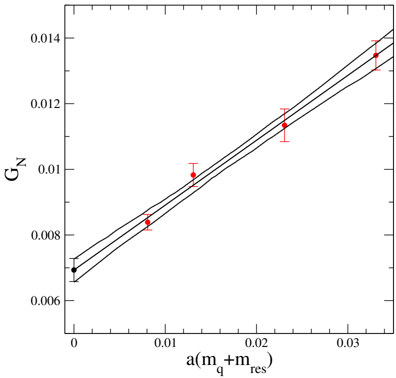
Nucleon Mass and amplitude

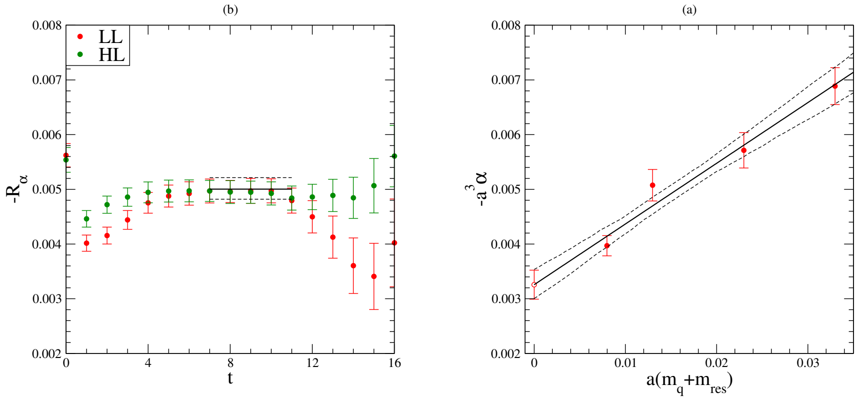


Nucleon Extrapolations

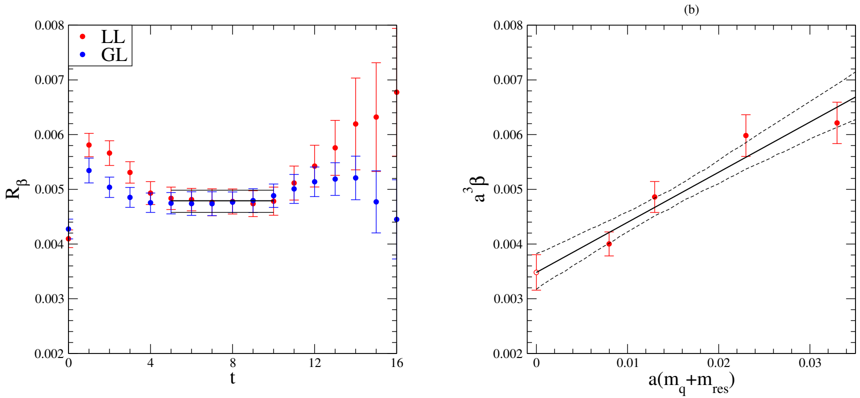




Low energy constant: α



Low energy constant: β



Systematic Errors

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- Chiral Extrapolation errors
- (Continuum Extrapolation errors)
- Errors in renormalisation
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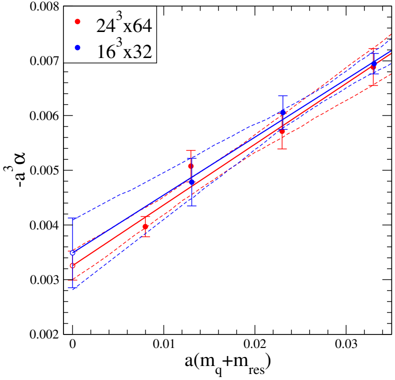
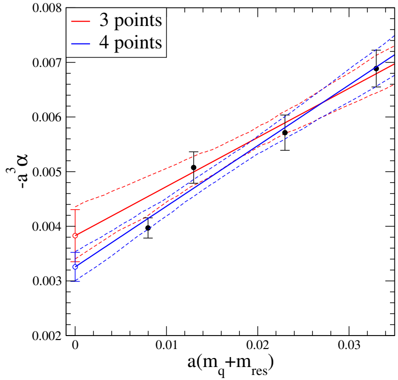
Finite Volume Error



- No noticeable effect

Extrapolation Error



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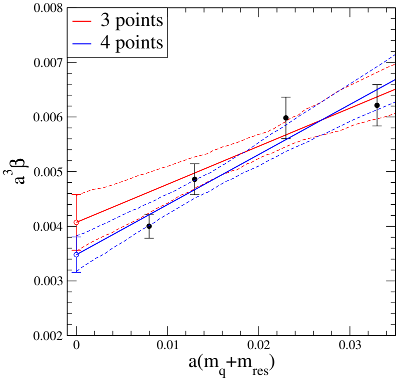

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NPR

- Non-perturbative MOM scheme renormalisation of the Rome-Southampton group
- The renormalised operators are

$$\mathcal{O}^{A}_{\mathrm{ren}} = Z^{AB}\mathcal{O}^{B}_{\mathrm{latt}}$$

- A and B label the spin structure, eg LL
- Z^{AB} is the mixing matrix
- $\mathcal{O}^{LL}$ and $\mathcal{O}^{RL}$ mix with a 3rd operator $\mathcal{O}^{A(LV)}$
  $\Rightarrow$ Z^{AB} is a 3×3 matrix
- Exponentially accurate chiral symmetry from Domain Wall Fermions should suppress operator mixing

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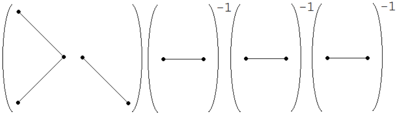
We want to calculate the non-perturbative amputated 3-quark vertex function of these operators



$$\mathcal{G}^{A}_{abc,\alpha\beta\gamma\delta}(p^2) = \epsilon^{abc}(C\Gamma)_{\alpha'\beta'}\Gamma'_{\delta\gamma'}\langle Q^{a'a}_{\alpha'\alpha}(p)Q^{b'b}_{\beta'\beta}(p)Q^{c'c}_{\gamma'\gamma}(p)\rangle$$

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and Γ and Γ' are the matrices which appear in $\mathcal{O}^A$

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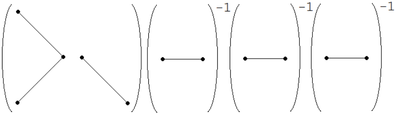

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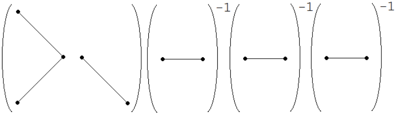

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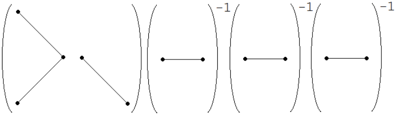

$$\mathcal{G}^{A}_{abc,\alpha\beta\gamma\delta}(p^2) = \epsilon^{abc}(C\Gamma)_{\alpha'\beta'}\Gamma'_{\delta\gamma'}\langle Q^{a'a}_{\alpha'\alpha}(p)Q^{b'b}_{\beta'\beta}(p)Q^{c'c}_{\gamma'\gamma}(p)\rangle$$

where

$$Q^{a'a}_{\alpha'\alpha} = \langle S^{a'a''}_{\alpha'\alpha''}(p)\rangle^{-1} S^{a''a}_{\alpha''\alpha}(p)$$

and Γ and Γ' are the matrices which appear in $\mathcal{O}^A$

- The renormalization condition in the RI–Mom Scheme is

$$Z_q^{-3/2} Z^{BC} M^{CA} = \delta^{BA}$$

- Where the matrix M is,

$$M^{AB} = \mathcal{G}^{A}_{abc,\alpha\beta\gamma\delta}(p^2) P^{B}_{abc,\beta\alpha\delta\gamma}$$

- and the projection matrices $P^{A}_{abc,\beta\alpha\delta\gamma}$ are chosen so that the renormalization condition is satisfied in the free field case where $Z_q = 1$ and $Z^{BC} = \delta^{BC}$.
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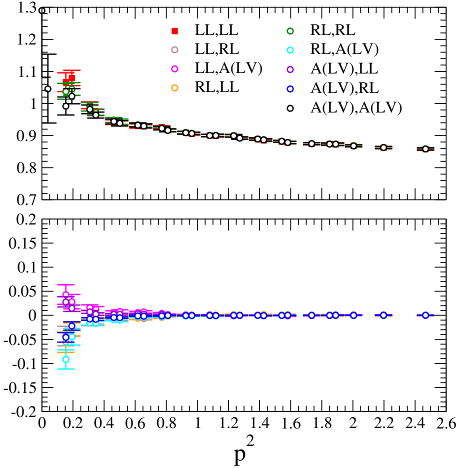

1.3
1.2
1.1
1
0.9
0.8
0.7
LL,LL
LL,RL
LL,A(LV)
RL,LL
RL,RL
RL,A(LV)
A(LV),LL
A(LV),RL
A(LV),A(LV)
0.2
0.15
0.1
0.05
0
-0.05
-0.1
-0.15
-0.2
0
0.2
0.4
0.6
0.8
1
1.2
1.4
1.6
1.8
2
2.2
2.4
2.6
p^2

- Perform a chiral extrapolation
- Match to the $\overline{\mathrm{MS}}$ scheme at 2GeV
- This gives
  $U^{\overline{\mathrm{MS}}\leftarrow \mathrm{latt}}(2GeV)_{LL} = 0.662(10)$
  $U^{\overline{\mathrm{MS}}\leftarrow \mathrm{latt}}(2GeV)_{RL} = 0.664(8)$
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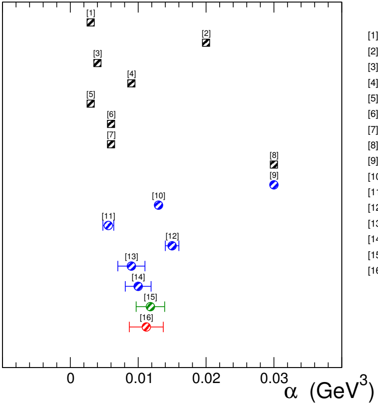
Summary



[1] Donoghue 82
[2] Thomas 83
[3] Meljanac 82
[4] Ioffe 81
[5] Krasnikov 82
[6] Ioffe 84
[7] Tomozawa 81
[8] Brodsky 84
[9] Hara 86
[10] Bowler 88
[11] Gavela 89
[12] JLQCD 00
[13] CP-PACS & JLQCD 04
[14] RBC 07
[15] RBC 07
[16] This work

- $\alpha = -0.0112(12)(22)$
- $\beta = 0.0120(13)(23)$

Outlook

- The direct calculation is currently underway
- Example: Preliminary results for the $W_0^{LL}(p \to \pi^+ + \nu)$, on the $16^3 \times 32$ lattice, with valence quark mass $am_u = 0.03$

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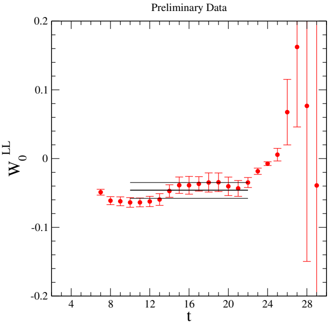

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- We can put bounds on the GUT scale physics:

Decay Mode	Lifetime bound(yrs)	A_{GUT} bound (M_{GUT}^{-4})
$p \longrightarrow e^+\pi^0$	$> 8.2 \times 10^{33}$	< 44
$p \longrightarrow e^+\pi^0$	$> 8.2 \times 10^{33}$	< 37
$p \longrightarrow K^+\bar{\nu}$	$> 2.3 \times 10^{33}$	< 76
$n \longrightarrow K^0\bar{\nu}$	$> 1.3 \times 10^{32}$	< 733

- For example:
  $p \rightarrow \pi^0 e^+$ via X Boson Exchange
  Minimal SU(5) SUSY GUT
- A_{GUT} is given by

$$A_{\mathrm{GUT}} = \frac{g_5^4 A_R^2}{M_X^4} \left|1 + (1 + V_{ud}^2)^2\right|$$

- Can put a bound on M_X
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